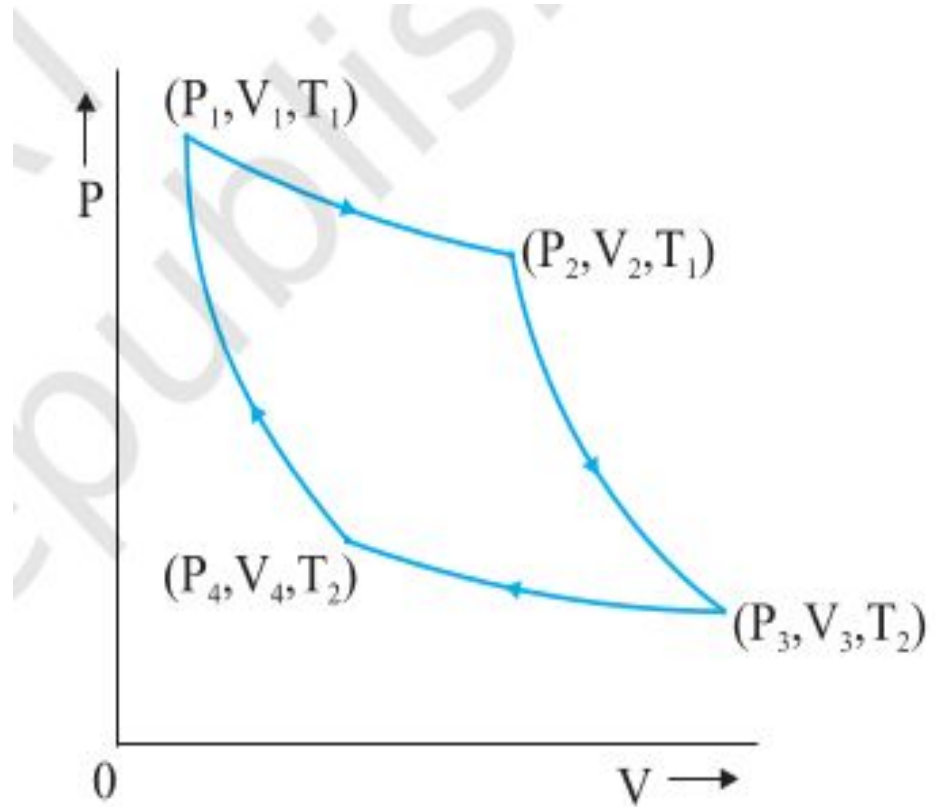


# CARNOT ENGINE

We expect the ideal engine operating between two temperatures to be a reversible engine

Suppose we have a hot reservoir at temperature  $T_1$  and a cold reservoir at temperature  $T_2$



A reversible heat engine operating between two temperatures is called a Carnot engine.

We have taken the working substance of the Carnot engine to be an ideal gas.

## Isothermal expansion

Isothermal expansion of the gas taking its state from

$$(P_1, V_1, T_1) \text{ to } (P_2, V_2, T_1).$$

The heat absorbed by the gas ( $Q_1$ ) from the reservoir at temperature  $T_1$

# Work done in Isothermal Expansion

Work done **by the gas** on the environment.

$$W_{1 \rightarrow 2} = Q_1 = nRT_1 \ln \left[ \frac{V_2}{V_1} \right]$$

# Adiabatic expansion

Adiabatic expansion of the gas from

$$(P_2, V_2, T_1) \text{ to } (P_3, V_3, T_2)$$

Work done **by the gas**

$$W_{2-3} = \left[ \frac{\mu R (T_1 - T_2)}{\gamma - 1} \right]$$

## Isothermal compression

Isothermal compression of the gas from

$$(P_3, V_3, T_2) \text{ to } (P_4, V_4, T_2).$$

Work done **on the Gas**

$$W_{3 \rightarrow 4} = Q_2 = \mu RT_2 \ln \left[ \frac{V_3}{V_4} \right]$$

# Adiabatic compression

Adiabatic compression of the gas from

$$(P_4, V_4, T_2) \text{ to } (P_1, V_1, T_1)$$

Work done **on the gas**

$$W_{4 \rightarrow 1} = \left[ \frac{\mu R (T_1 - T_2)}{\gamma - 1} \right]$$



Total work done

Total work done by the gas in one complete cycle is given by

$$W = W_{1 \rightarrow 2} + W_{2 \rightarrow 3} - W_{3 \rightarrow 4} - W_{4 \rightarrow 1}$$

$$W = \mu RT_2 \ln \left[ \frac{V_3}{V_4} \right] - \mu RT_2 \ln \left[ \frac{V_3}{V_4} \right]$$

Efficiency

$$\eta = 1 - \frac{Q_2}{Q_1}$$

And

$$\eta = \frac{W}{Q_1} = \frac{\mu RT_1 \ln \left[ \frac{V_3}{V_4} \right] - \mu RT_2 \ln \left[ \frac{V_3}{V_4} \right]}{\mu RT_1 \ln \left[ \frac{V_3}{V_4} \right]}$$

On simplification

$$\eta = 1 - \frac{T_2 \ln \left[ \frac{V_3}{V_4} \right]}{T_1 \ln \left[ \frac{V_2}{V_1} \right]}$$

We can find that

$$\left[ \frac{V_3}{V_4} \right] = \left[ \frac{V_2}{V_1} \right] \text{ In the following way}$$

In the Adiabatic Expansion Process ( $P_2, V_2, T_1$ ) to ( $P_3, V_3, T_2$ )

$$P_2 V_2^\gamma = P_3 V_3^\gamma$$

But  $P_2 V_2 = \mu RT_1$  and  $P_3 V_3 = \mu RT_2$

$$\mu RT_1 V_2^{\gamma-1} = \mu RT_2 V_3^{\gamma-1}$$

$$T_1 V_2^{\gamma-1} = T_2 V_3^{\gamma-1}$$

Cross Multiplying  $\left[ \frac{V_2}{V_3} \right] = \left[ \frac{T_2}{T_1} \right]^{\frac{1}{(\gamma-1)}}$

Similarly in the adiabatic compression process

$(P_4, V_4, T_2)$  to  $(P_1, V_1, T_1)$

$$\left[ \frac{V_1}{V_4} \right] = \left[ \frac{T_2}{T_1} \right]^{\frac{1}{(\gamma-1)}}$$

## On Comparison

$$\left[ \frac{V_2}{V_3} \right] = \left[ \frac{V_1}{V_4} \right] \quad \text{Or} \quad \left[ \frac{V_2}{V_1} \right] = \left[ \frac{V_3}{V_4} \right]$$

Then we prove that

$$\eta = 1 - \frac{T_2}{T_1}$$

In Carnot Cycle

$$\frac{T_2}{T_1} = \frac{Q_2}{Q_1}$$

It is a universal relation independent of the nature of the system